



Generalised hydrodynamics of the KdV soliton gas

Dynamics Days Europe
MS: Recent advances in dispersive and
generalised hydrodynamics

Thibault Bonnemain, 8th September 2023 [joint work with Benjamin Doyon and Gennady El]

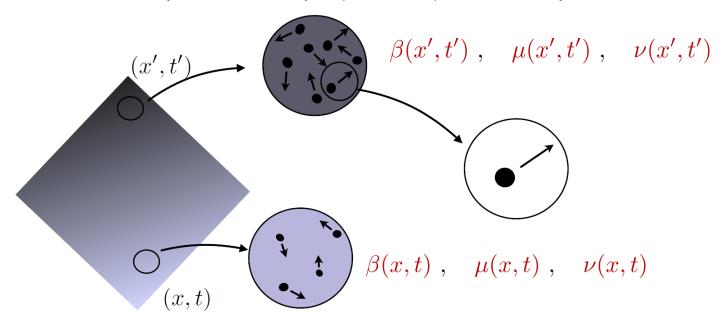
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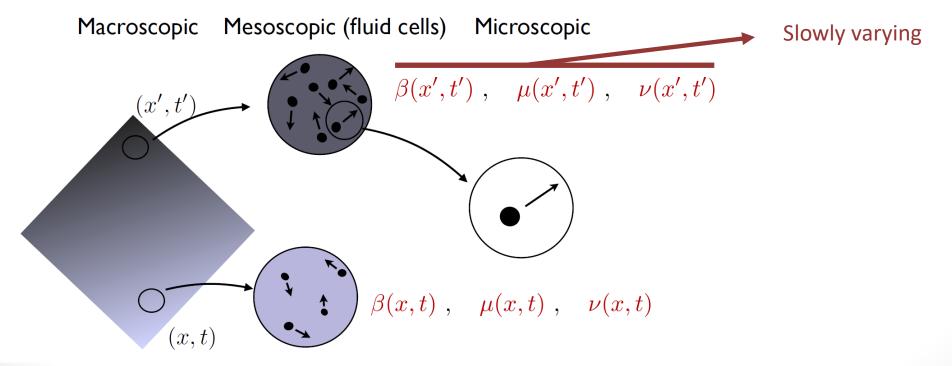
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- Hydrodynamic principle: separation of scales and propagation of local GE

Macroscopic Mesoscopic (fluid cells) Microscopic

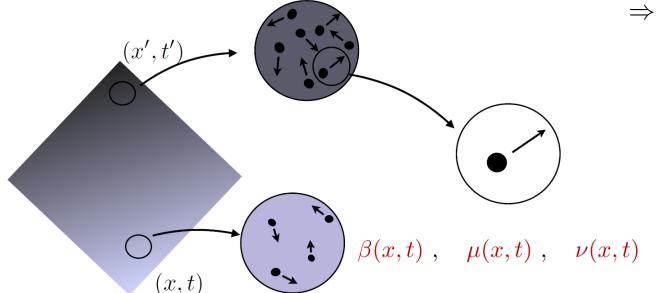


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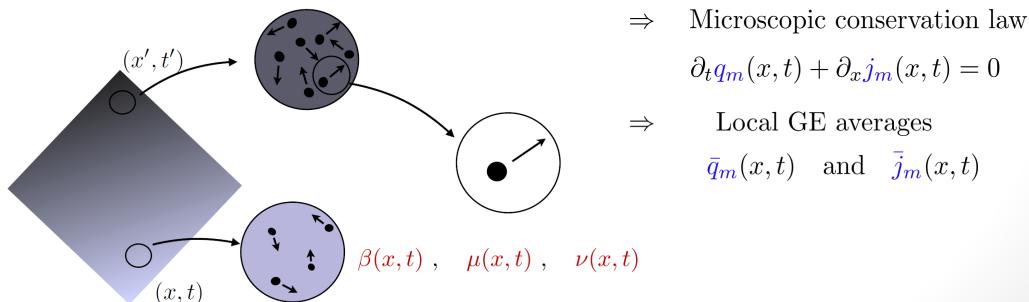


Microscopic conservation law

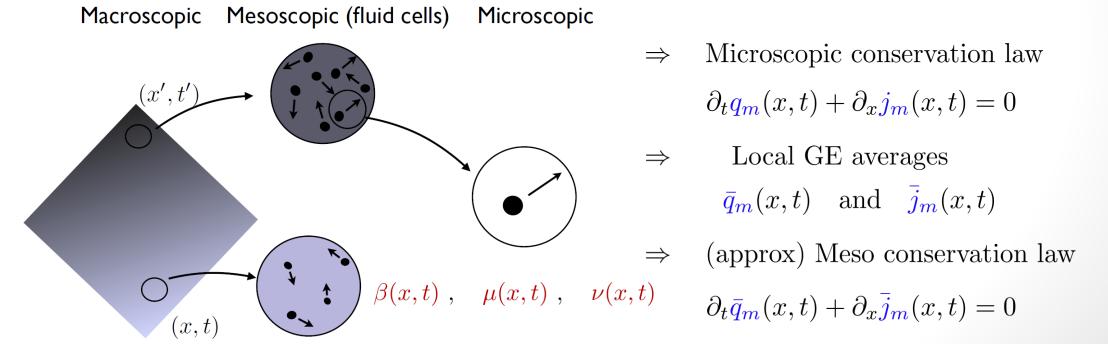
$$\partial_t \mathbf{q}_m(x,t) + \partial_x \mathbf{j}_m(x,t) = 0$$

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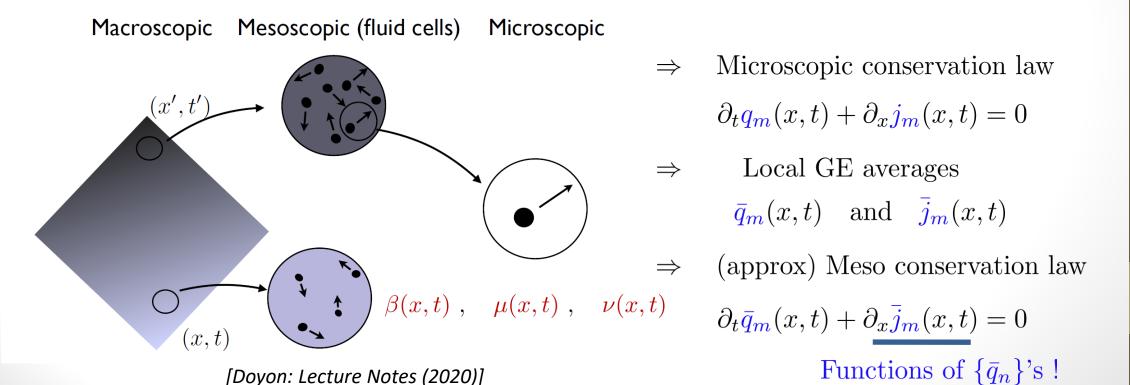
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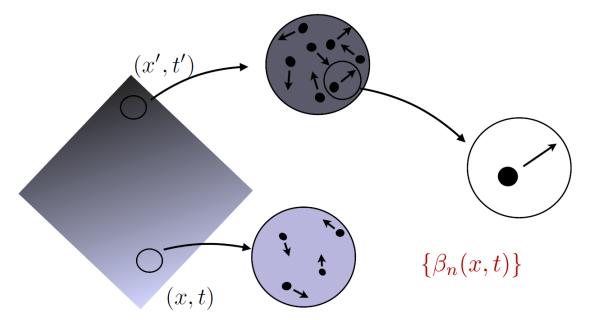
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Generalised Gibbs ensembles and GHD

- Boltzmann 1868: micro-canonical ensemble in long time limit
 - \triangleq Generalised Gibbs ensembles (GE): $\rho \propto e^{-\sum_{n=0}^{\infty} \beta_n Q_n}$
- Hydrodynamic principle: separation of scales and propagation of local GGE

Macroscopic Mesoscopic (fluid cells) Microscopic



[Doyon: Lecture Notes (2020)]

 \Rightarrow Microscopic conservation law

$$\partial_t \mathbf{q_m}(x,t) + \partial_x \mathbf{j_m}(x,t) = 0$$

 \Rightarrow Local GE averages

$$\bar{q}_m(x,t)$$
 and $\bar{j}_m(x,t)$

 \Rightarrow (approx) Meso conservation law

$$\partial_t \bar{q}_m(x,t) + \partial_x \bar{j}_m(x,t) = 0$$

Functions of $\{\bar{q}_n\}$'s!

Korteweg - De Vries equation

• KdV: integrable nonlinear dispersive PDE

$$\partial_t \varphi + 6\varphi \partial_x \varphi + \partial_x^3 \varphi = 0 .$$

Korteweg - De Vries equation

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• Infinite set of conservation laws

Time conserved "charges"
$$Q_n = \int \mathrm{d}x \; q_n(x,t) \;, \quad J_n = \int \mathrm{d}t \; j_n(x,t) \;, \qquad \text{Space conserved conserved "currents"}$$

$$\partial_t q_n + \partial_x j_n = 0 .$$

[Miura, Gardner, Kruskal (1968)]

Korteweg - De Vries equation

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$$\partial_t q_n + \partial_x j_n = 0 .$$

• Solvable via Inverse Scattering Transfom.

• KdV: integrable nonlinear dispersive PDE

$$\partial_{t}\varphi + 6\varphi\partial_{x}\varphi + \partial_{x}^{3}\varphi = 0.$$

$$1.8$$

$$1.4$$

$$1.2$$

$$1$$

$$0.8$$

$$0.6$$

$$0.4$$

$$0.2$$

$$-250$$

$$-200$$

$$-150$$

$$-100$$

$$-50$$

$$0$$

$$50$$

$$100$$

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$$250$$

Single realisation of a KdV soliton gas

0.4

0.2

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 Single realisation of a KdV soliton gas

50

0 x 100

150

200

250

Fluid cell of size L characterised by local GGE

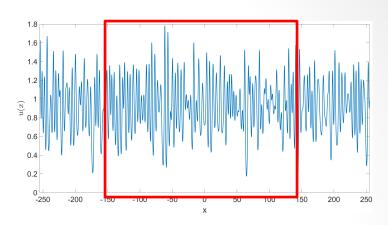
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• Multisoliton solution

$$\varphi_N \sim \sum_{i=1}^N 2\eta_i^2 \operatorname{sech}^2\left[\eta_i \left(x - 4\eta_i^2 t - x_i^{\pm}\right)\right] \text{ as } t \to \pm \infty.$$

N solitons in L



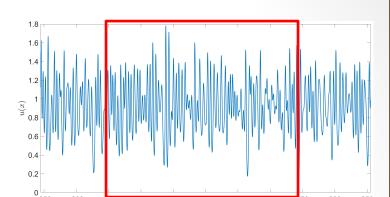
as
$$t \to \pm \infty$$
.

[Zakharov (1971)]

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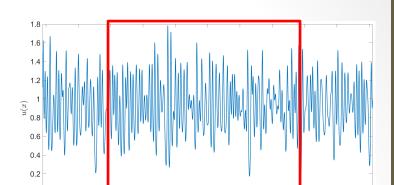
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 Angle coordinate

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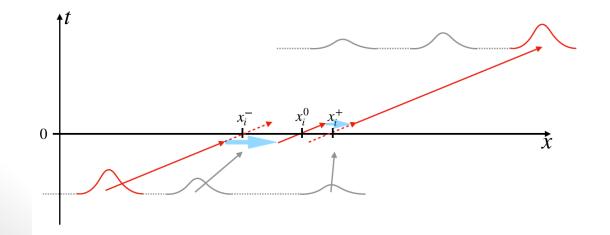
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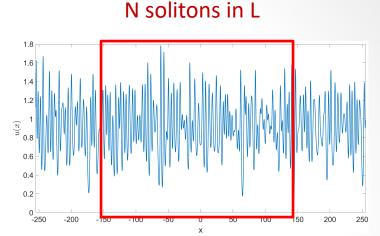


Scattering is elastic and 2-body factorisable

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 Angle coordinate

• Relation between asymptotic states given by scattering shift

$$\left| \begin{array}{c} x_i^+ - x_i^- = \sum_i rac{ ext{sgn}(\eta_{ ext{i}} - \eta_{ ext{j}})}{\eta_i} \ln \left| rac{\eta_i + \eta_j}{\eta_i - \eta_j}
ight| \end{array}
ight|$$

[Lax (1968)]

• Partition function

$$\mathcal{Z} = \sum_{N=0}^{\infty} \frac{1}{N!} \int_{\Gamma^N \times \mathbb{R}^N} \prod_{i=1}^N \frac{\mathrm{d}p(\eta_i)}{2\pi} \mathrm{d}x_i^- \exp\left[-\sum_{i=1}^N w(\eta_i)\right] \chi\left(\varphi_N(x, t=0) < \epsilon_x, x \notin [0, L]\right)$$

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Soliton bare velocity

$$p(\eta) = 4\eta^2$$

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Generalised
Gibbs weights

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 e.g. $w(\eta) = \sum_{k} \beta_k h_k(\eta)$

 $h_n(\eta) = Q_n$ for a single soliton η

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Soliton bare velocity

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Constraint / Entropy

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• Thermodynamic limit $L \to \infty$

$$\mathcal{Z} symp \exp\left(-L\mathcal{F}\right) \; , \qquad \mathcal{F} = -\int_{\Gamma} \frac{\eta \mathrm{d}\eta}{\sigma(\eta)}$$

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 Spectral scaling function

• NDR of soliton gases

$$\sigma(\eta)\rho(\eta) = \eta - \int_{\Gamma} d\mu \, \rho(\mu) \log \left| \frac{\eta + \mu}{\eta - \mu} \right|$$

$$ho(\eta) d\eta dx = \# ext{ of solitons in } [x, x + dx] \times [\eta, \eta + d\eta]$$
 sity of States

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• Alternative interpretation

$$\frac{\mathrm{d}x^{-}(\eta)}{\mathrm{d}x} = \frac{\sigma(\eta)\rho(\eta)}{\eta}$$

Change of metric

• Spectral scaling function

$$\log \left[\frac{4\sigma(\eta)}{\pi} \right] = w(\eta) - \int_{\Gamma} d\mu \frac{1}{\sigma(\mu)} \log \left| \frac{\eta - \mu}{\eta + \mu} \right|, \qquad \mathcal{F} = -\int_{\Gamma} \frac{\eta d\eta}{\sigma(\eta)}$$

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• Density of states

$$\rho(\eta) = \frac{\delta \mathcal{F}}{\delta w(\eta)}$$

• Thermodynamic averages

$$\langle q_n \rangle = \int_{\Gamma} d\eta \ \rho(\eta) h_n(\eta) = \frac{\partial \mathcal{F}}{\partial \beta_n}$$

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$$= \int_{\Gamma} d\eta \, \rho(\eta) \theta(\eta) h_a^{\mathrm{dr}}(\eta) h_b^{\mathrm{dr}}(\eta)$$
[Doyon (2018)]

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$$\log \left[\frac{4\sigma(\eta)}{\pi} \right] = w(\eta) - \int_{\Gamma} d\mu \frac{1}{\sigma(\mu)} \log \left| \frac{\eta - \mu}{\eta + \mu} \right|, \qquad \mathcal{F} = -\int_{\Gamma} \frac{\eta d\eta}{\sigma(\eta)}$$

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$$\begin{aligned} \mathsf{C}_{ab} &\equiv \int_{\Gamma} \mathrm{d}x \left(\langle q_a(x) q_b(0) \rangle - \langle q_a(x) \rangle \langle q_b(0) \rangle \right) = -\frac{\partial^2 \mathcal{F}}{\partial \beta_a \partial \beta_b} \\ &= \int_{\Gamma} \mathrm{d}\eta \, \rho(\eta) \theta(\eta) h_a^{\mathrm{dr}}(\eta) h_b^{\mathrm{dr}}(\eta) & \quad \text{[Doyon (2018)]} \\ \text{Statistical factor} & h^{\mathrm{dr}}(\eta) = h(\eta) + \int_{\Gamma} \frac{\mathrm{d}\mu}{\sigma(\mu)} \log \left| \frac{\eta - \mu}{\eta + \mu} \right| h^{\mathrm{dr}}(\mu) \end{aligned}$$

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	MB	FD	BE	Simulations
C_{00}^{DC}	0.0235	-2.28	2.32	0.022 ± 0.003
C_{01}^{DC}	0.027	-3.18	3.23	0.024 ± 0.004
$egin{array}{c} {\sf C}^{ m DC}_{00} \ {\sf C}^{ m DC}_{01} \ {\sf C}^{ m DC}_{11} \ \end{array}$	0.042	-4.48	4.56	0.039 ± 0.005
C_{00}^{U}	0.22	0.028	0.41	0.2 ± 0.03
C^{U}_{01}	0.28	0.072	0.49	0.23 ± 0.04
$egin{array}{c} C_{00}^{\mathrm{U}} \ C_{01}^{\mathrm{U}} \ C_{11}^{\mathrm{U}} \end{array}$	0.39	0.12	0.66	0.36 ± 0.05

From thermodynamics to hydrodynamics

• Integrability: infinite number of conservation laws

$$\partial_t q_n + \partial_x j_n = 0$$

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• Hydrodynamic approximation: separation of scales

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$$\partial_{t} \rho(\eta; x, t) + \partial_{x} \left[v^{\text{eff}}(\eta; x, t) \rho(\eta; x, t) \right] = 0$$

$$v^{\text{eff}}(\eta) = 4\eta^{2} + \frac{1}{\eta} \int_{\Gamma} \log \left| \frac{\eta + \mu}{\eta - \mu} \right| \rho(\mu) \left[v^{\text{eff}}(\eta) - v^{\text{eff}}(\mu) \right] d\mu$$

Gibbs weights e.g.
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 Yang-Yang Equation
$$\sigma(\eta)$$

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$$w(\eta) = \sum_k \beta_k \eta^{2k+1}$$
 Yang-Yang Equation
$$\rho(\eta) \xrightarrow{\text{1st NDR}} \sigma(\eta)$$

Gibbs weights e.g.
$$w(\eta) = \sum_k \beta_k \eta^{2k+1}$$
 Yang-Yang Equation
$$\rho(\eta) \xrightarrow{\text{1st NDR}} \sigma(\eta)$$

$$v^{\text{eff}}(\eta) \xrightarrow{\text{ractric change}} v^{\text{eff}}(\eta)$$

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Yang-Yang Equation

$$\rho(\eta) \xrightarrow{\text{1st NDR}} \sigma(\eta)$$

$$v^{\text{eff}}(\eta) \xrightarrow{\text{nuctric change}} S$$

