

# A Mean-Field Game theoretical approach to crowd dynamics

University of Bristol

Thibault Bonnemain, 18th October 2023

*[Based on works w/ D. Ullmo, T. Gobron, M. Butano, C. Appert-Rolland, I. Echeverria-Huarte A. Nicolas, A. Seguin]*

# Pedestrian behavioral levels

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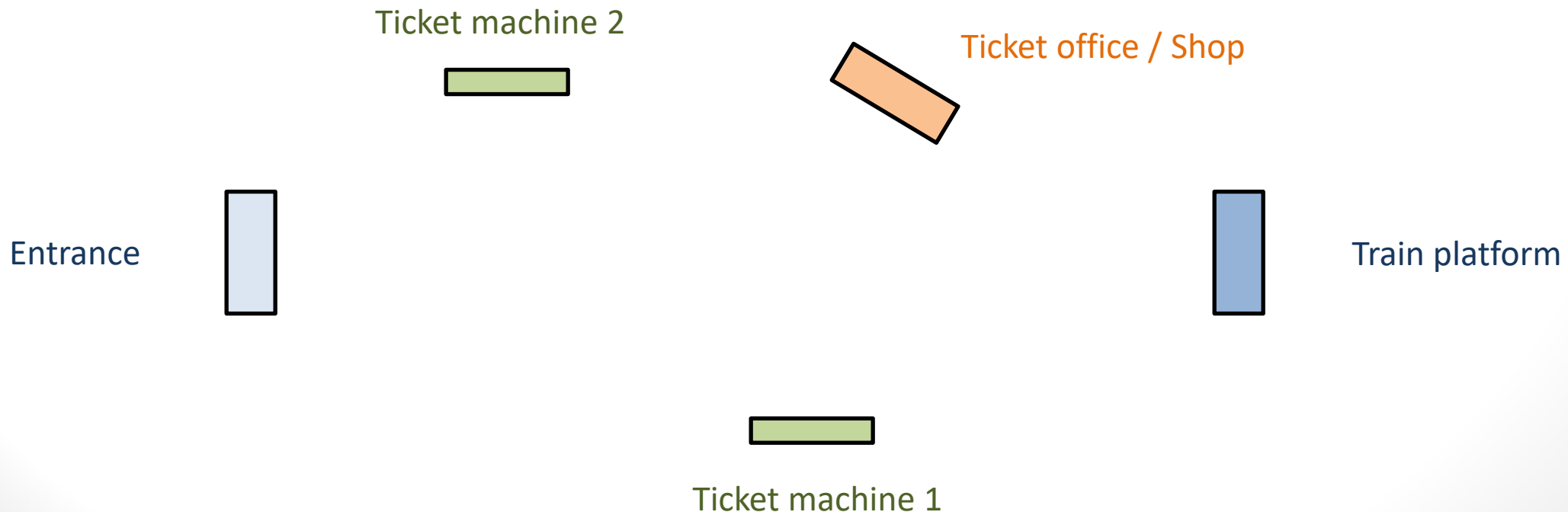
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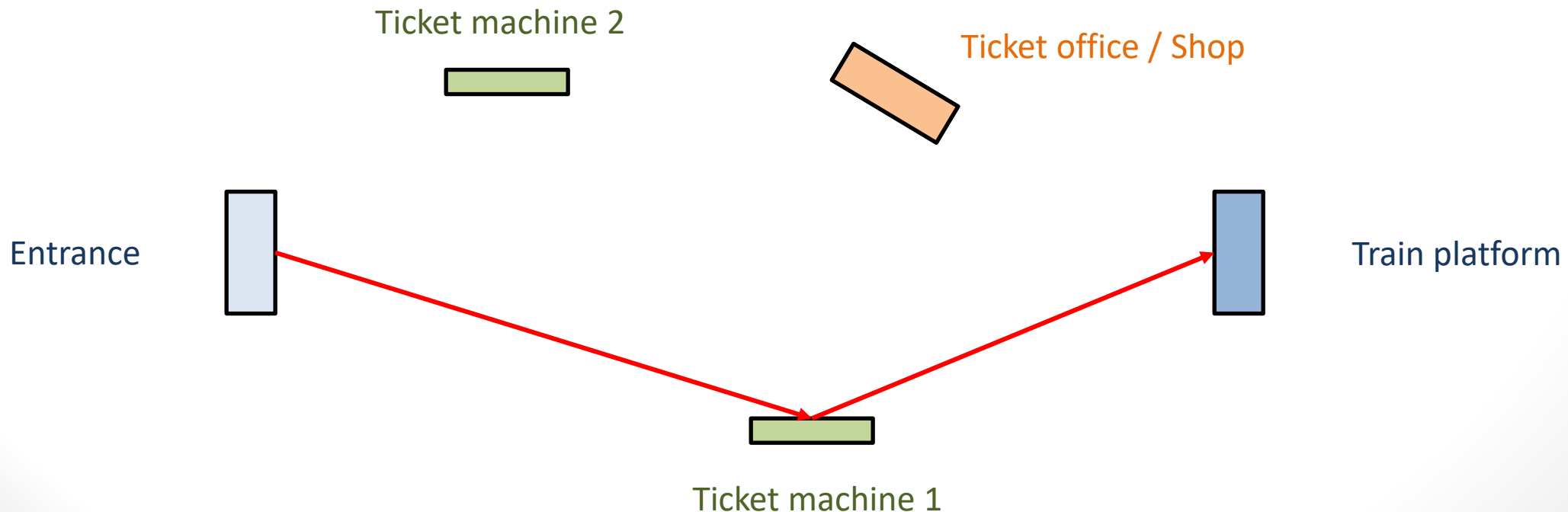
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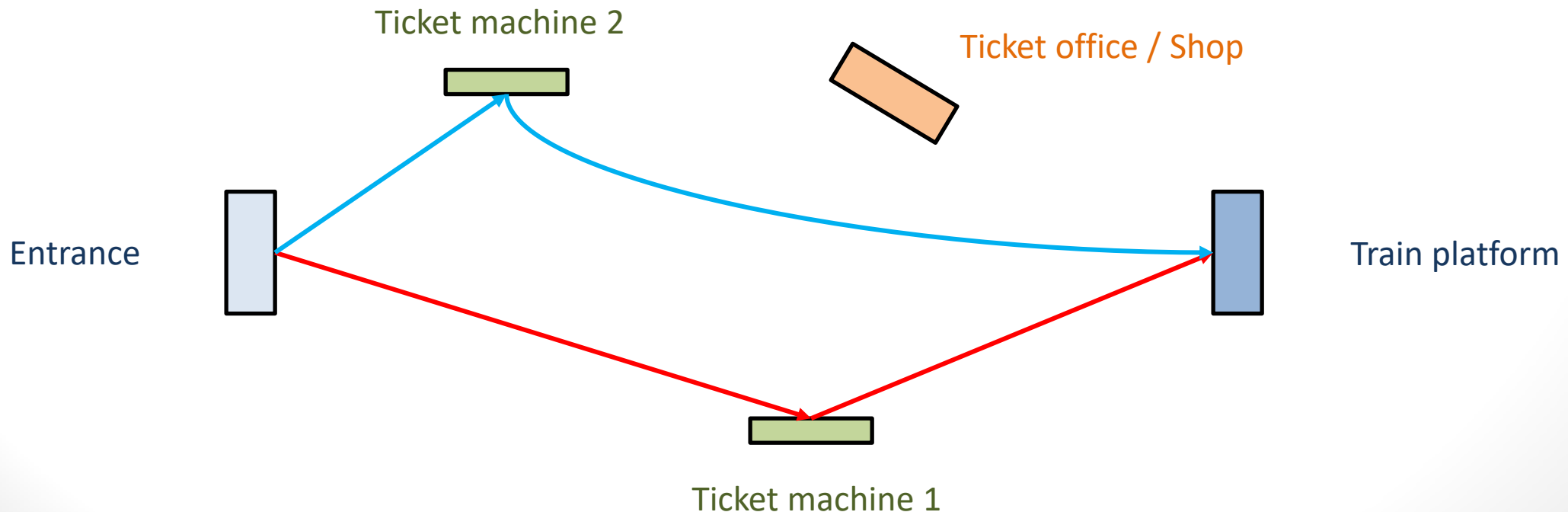
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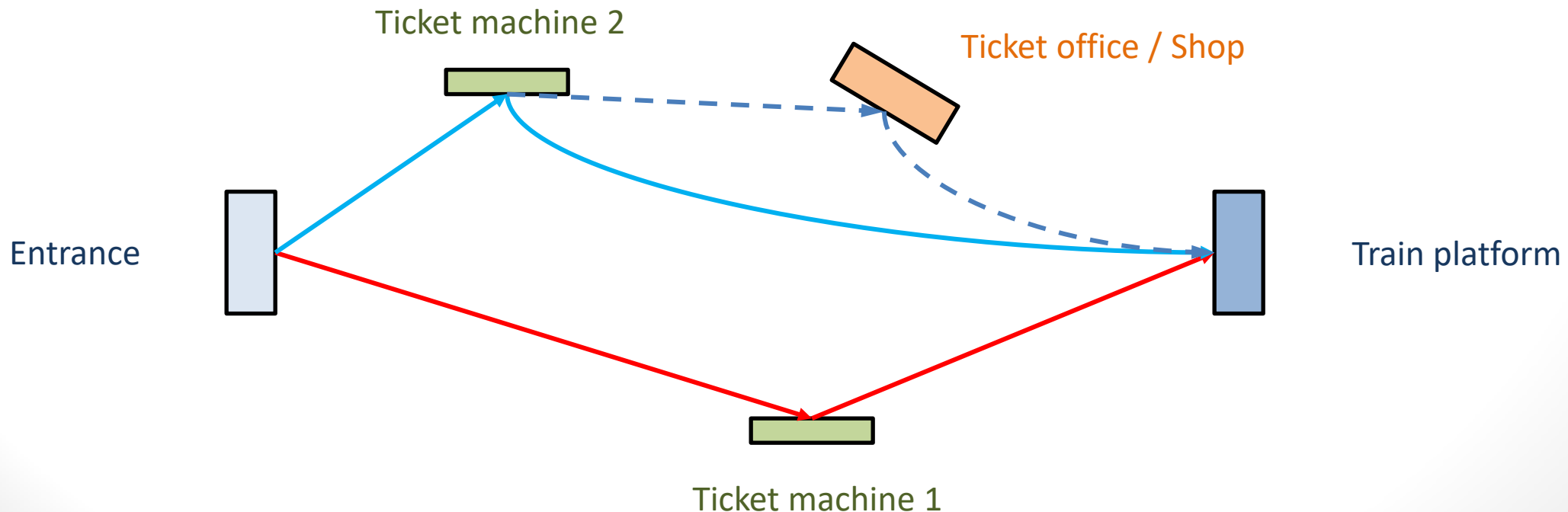




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# Intruder crossing a static crowd



(10)

(3)

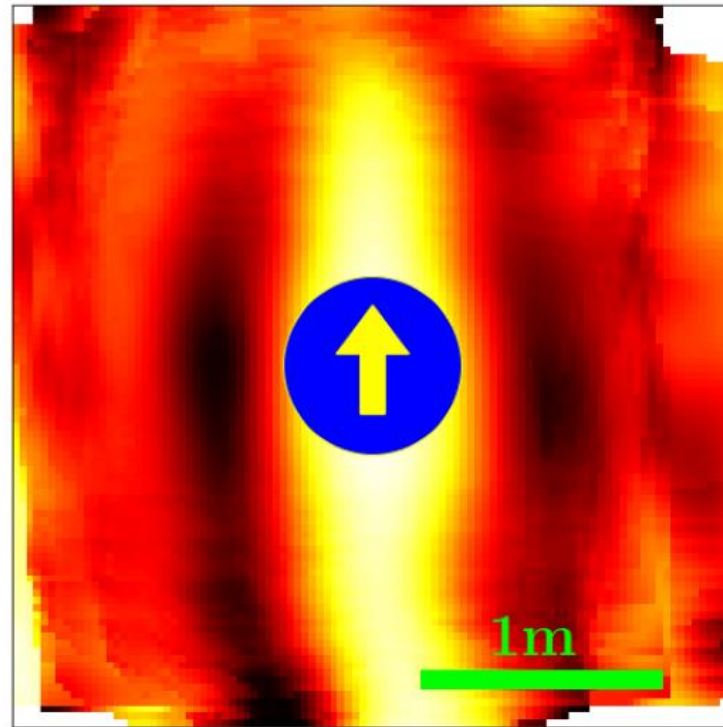
[Nicolas et al. (2019)]

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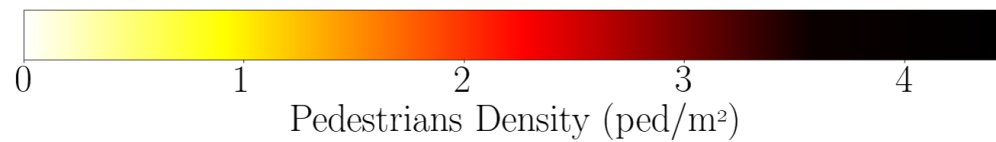
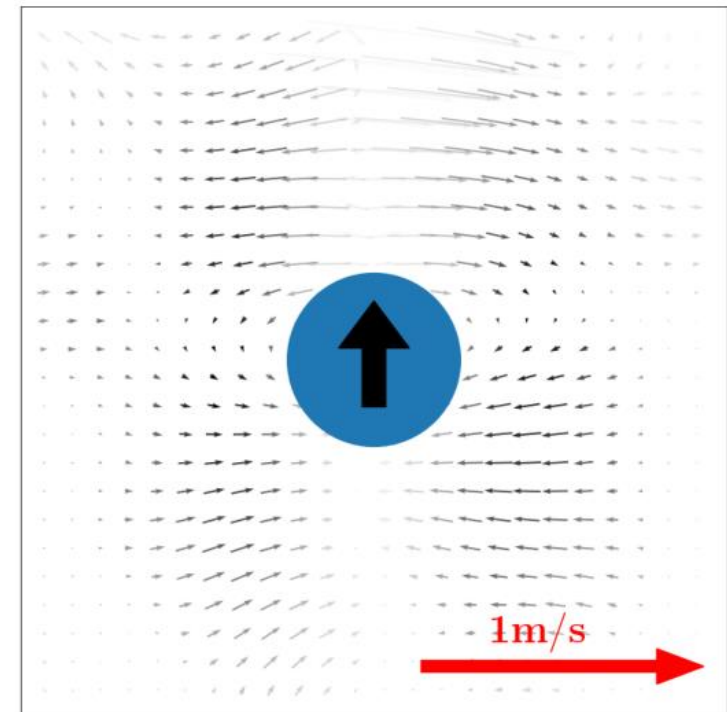
Experiment



(Averaged) pedestrian density



Velocity field



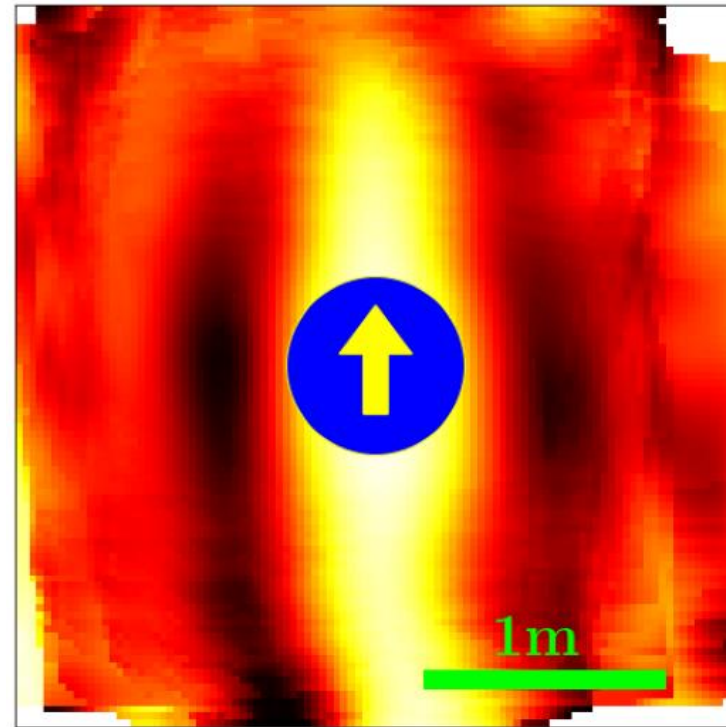


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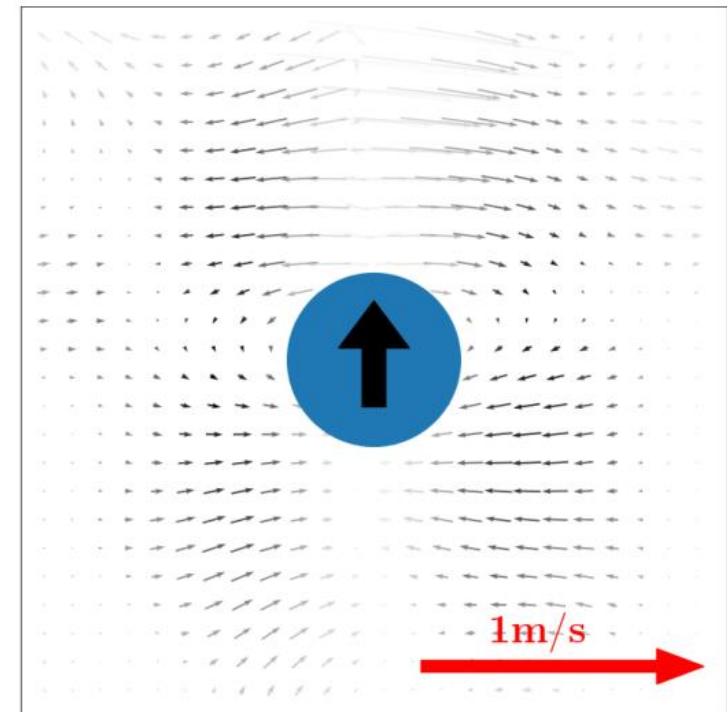
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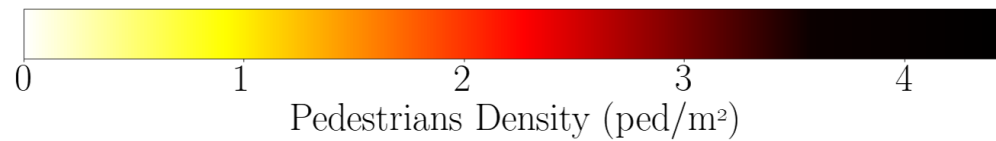
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Velocity field



Symmetric  
density profile

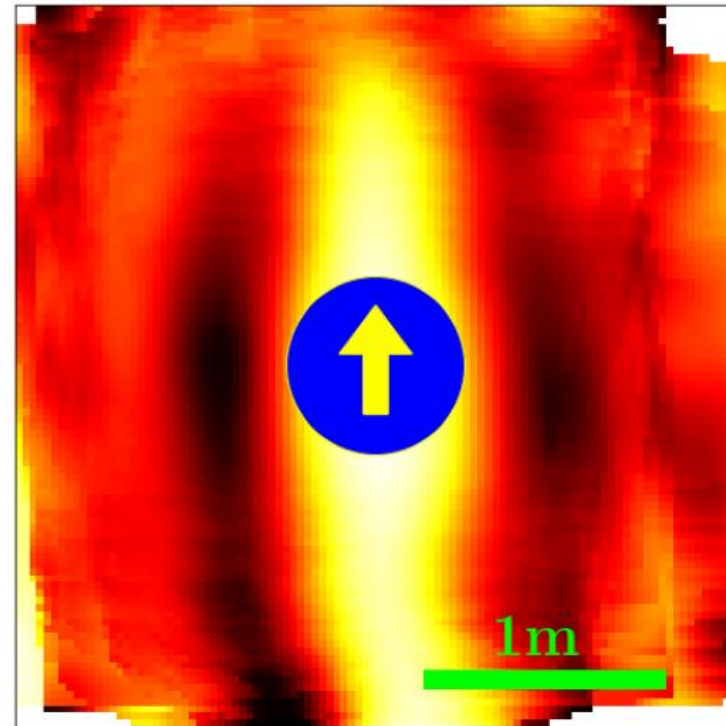


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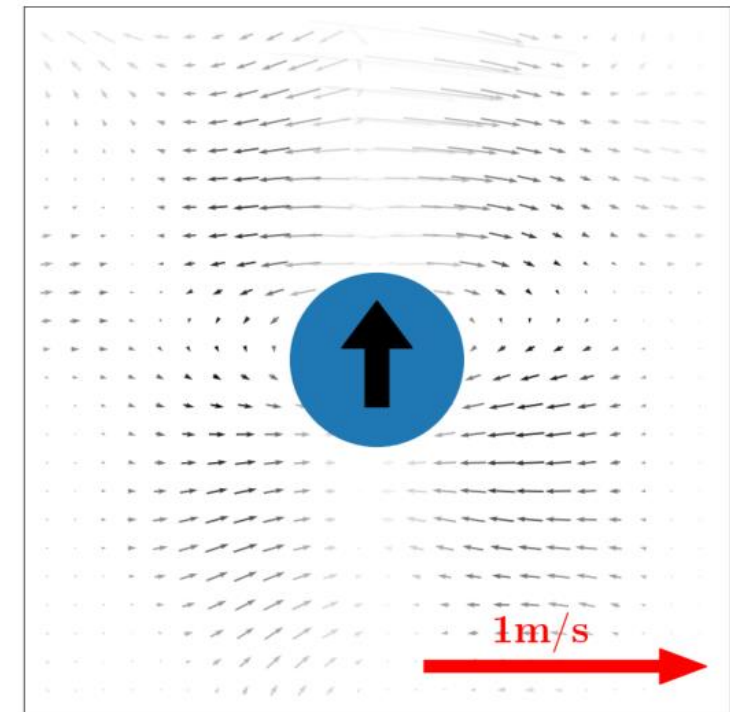
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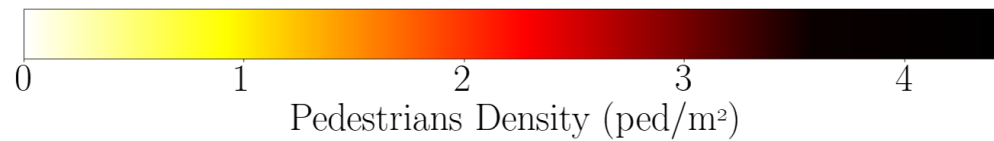
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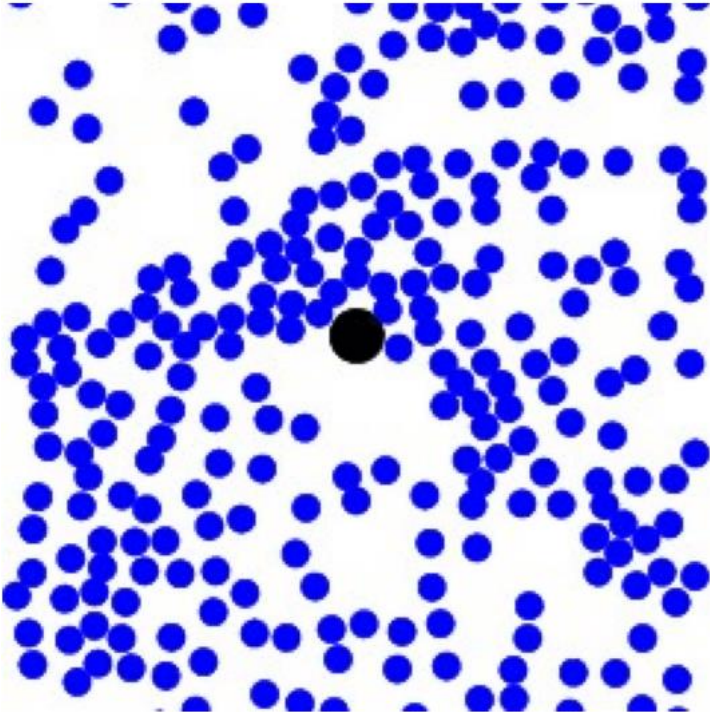
Transverse  
displacement



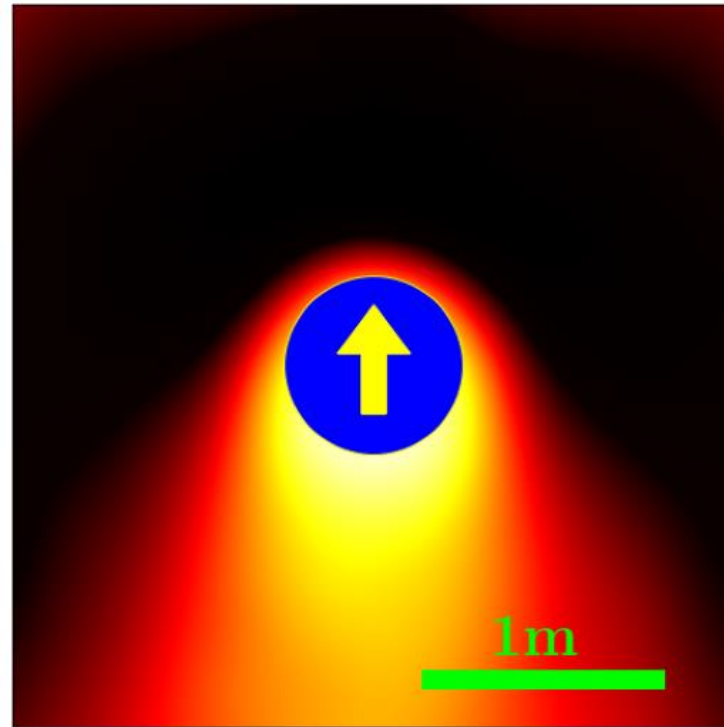
# Granular matter

Based on [Seguin et al. (2009)]

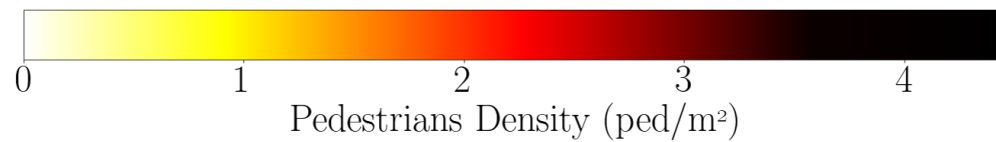
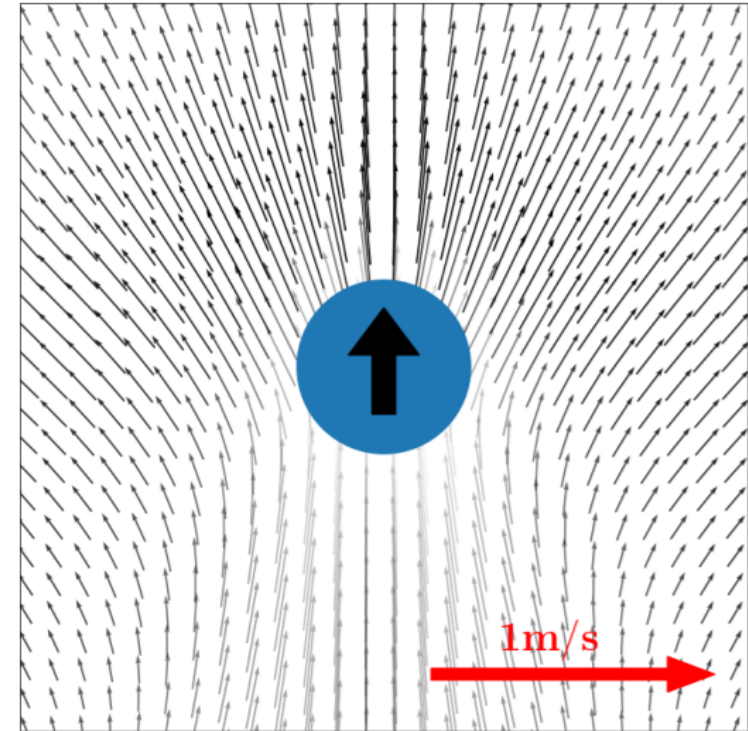
Simulation: monolayer of vibrated discs



(Averaged) pedestrian density



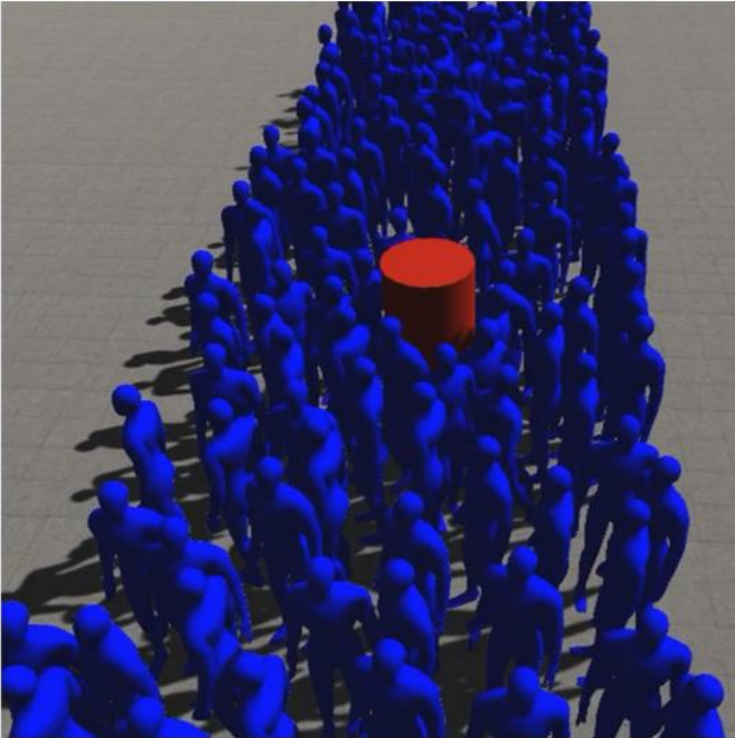
Velocity field



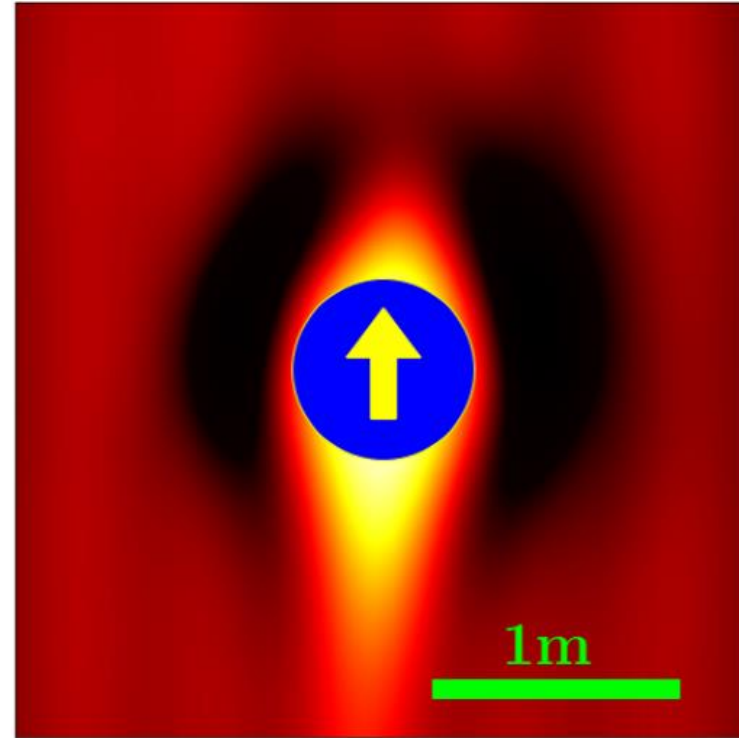
# Time To Collision models

Based on [Echeverria-Huarte, Nicolas(2023)]  
and [Karamouzas et al. (2017)]

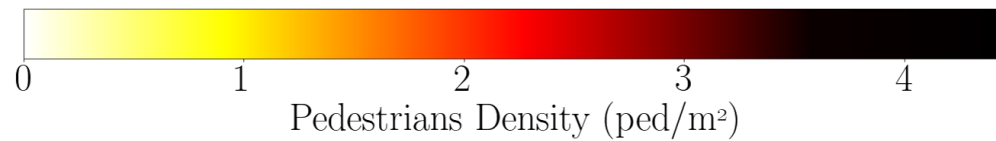
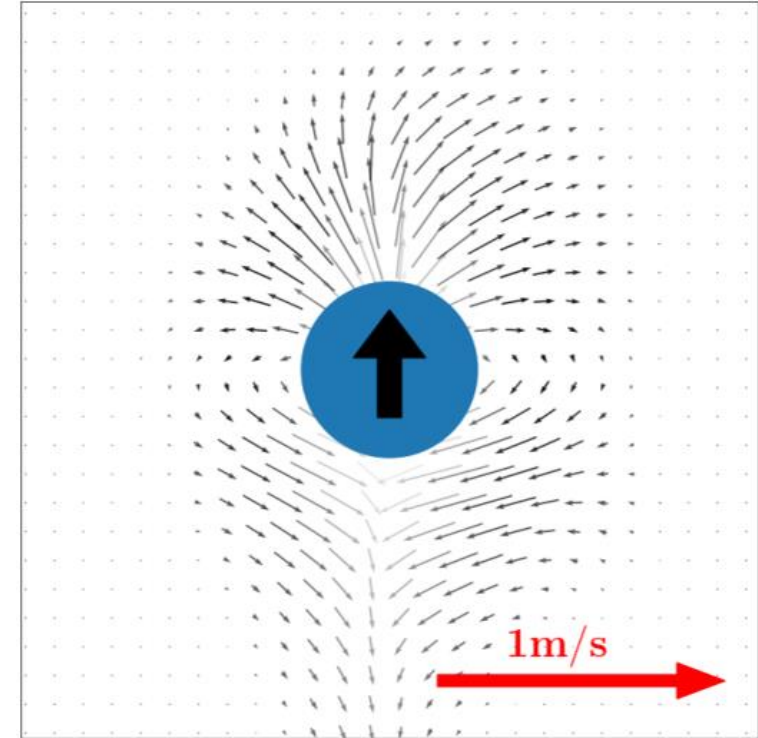
Simulation



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Velocity field



## To sum up...

- “Mechanical” models (granular and social forces) drastically fail to reproduce the qualitative features of the experiment.



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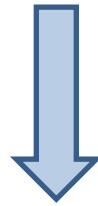
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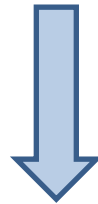


Change of paradigm

- Long-term anticipation → Competitive optimisation → Game theory

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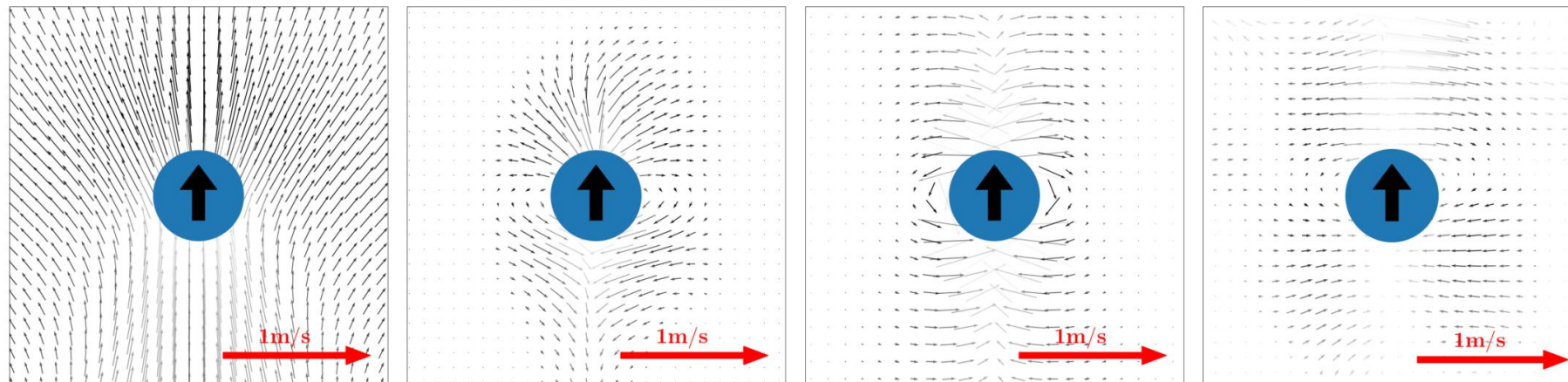
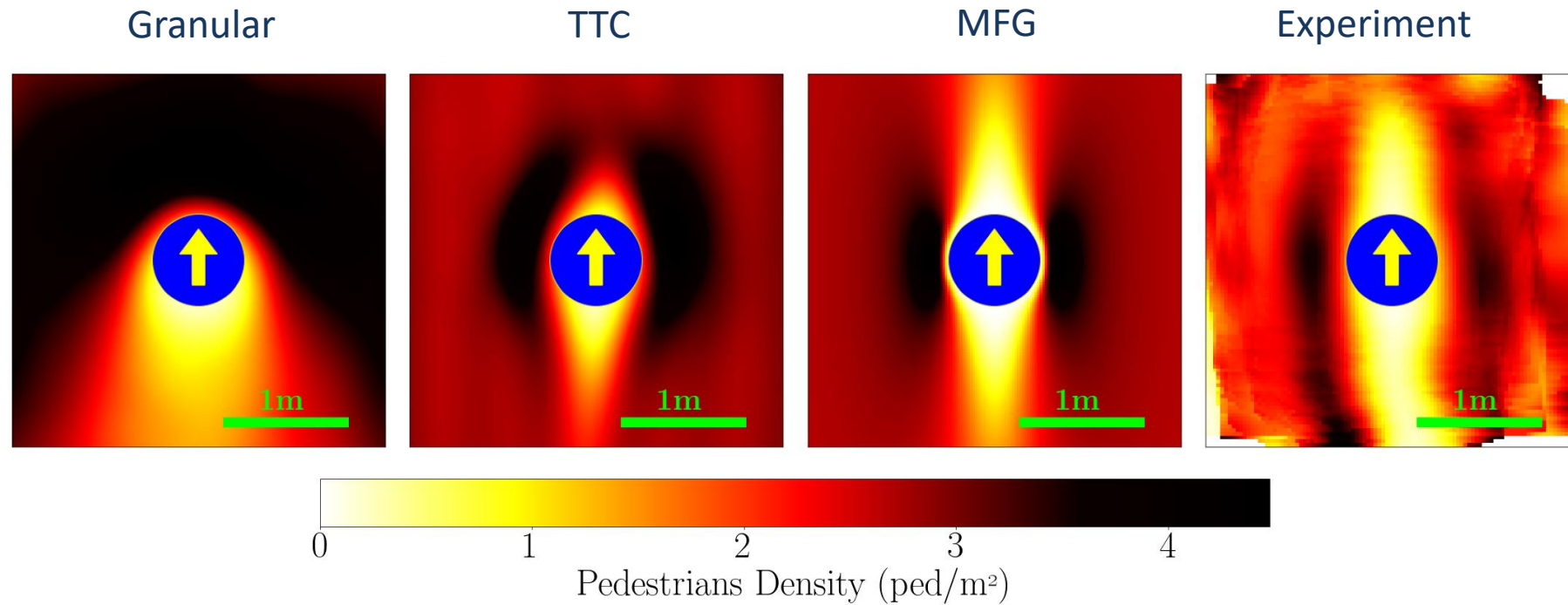


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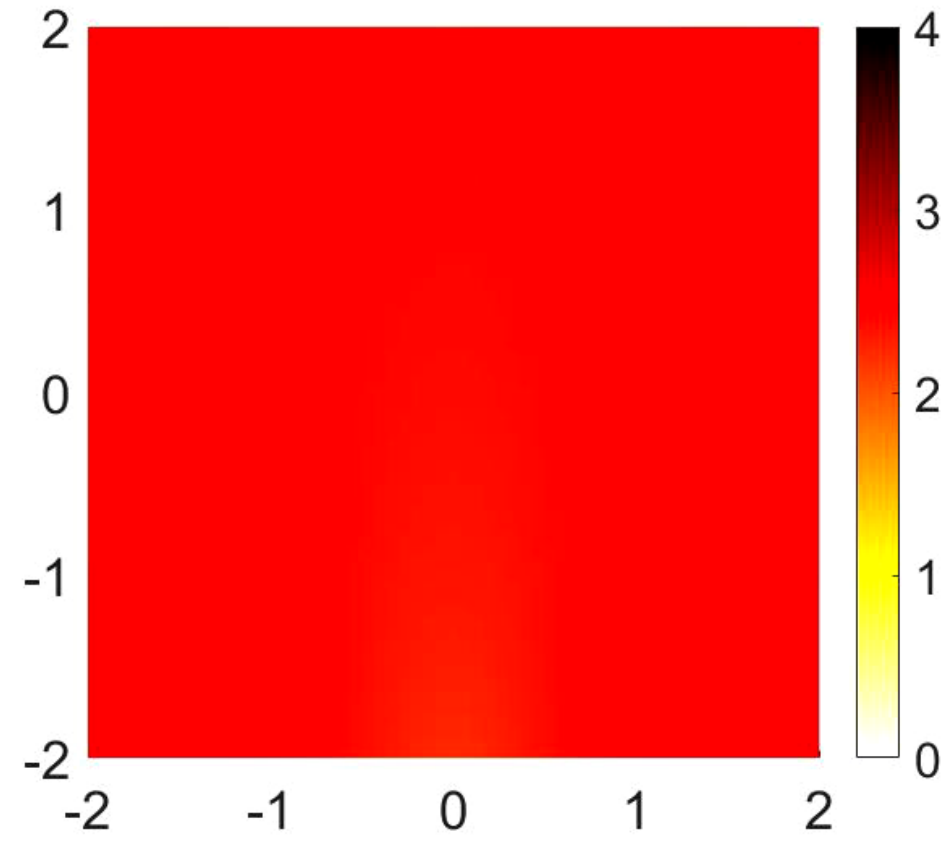
- Long-term anticipation → Competitive optimisation → Game theory
- Dense crowd → Many-body problem → Mean-field

# Comparing the different approaches

[Bonnemain et al. (2023)]



# Another example in a non-controlled environment



# Game theory

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Hawk-Dove paradigm

	H	D
H	$(\frac{1}{4}; \frac{1}{4})$	$(1; 0)$
D	$(0; 1)$	$(\frac{1}{2}; \frac{1}{2})$



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- Wide literature: mathematics, engineering sciences, economics, sociology ...

*[Lachapelle, Wolfram (2011)]*

*[Guéant et al. (2012)]*

*[Gomes, Saude (2014)]*

*[Laguzet, Turinici (2015)]*

*[Achdou et al. (2016)]*

*[Cardaliaguet, Lehalle (2017)]*

*[Bremaud, Ullmo (2022)]*

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# Mean Field Games equations

- Optimization: linear programming leads to a (backward) Hamilton-Jacobi-Bellman equation for the value function  $u(\mathbf{x}, t)$

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optimal control  $\mathbf{a}$

**Mean Field Game** = coupling between a (collective) stochastic motion and an (individual) optimization problem through a mean field

# Ergodic state

Theorem:

[Cardaliaguet, Lasry, Lions, Poretta (2013)]

- $V[\mathbf{m}](\mathbf{x}, t)$  has no explicit time dependence
- Long optimization time:  $T \rightarrow \infty$
- System is confined
- ... + other conditions ...



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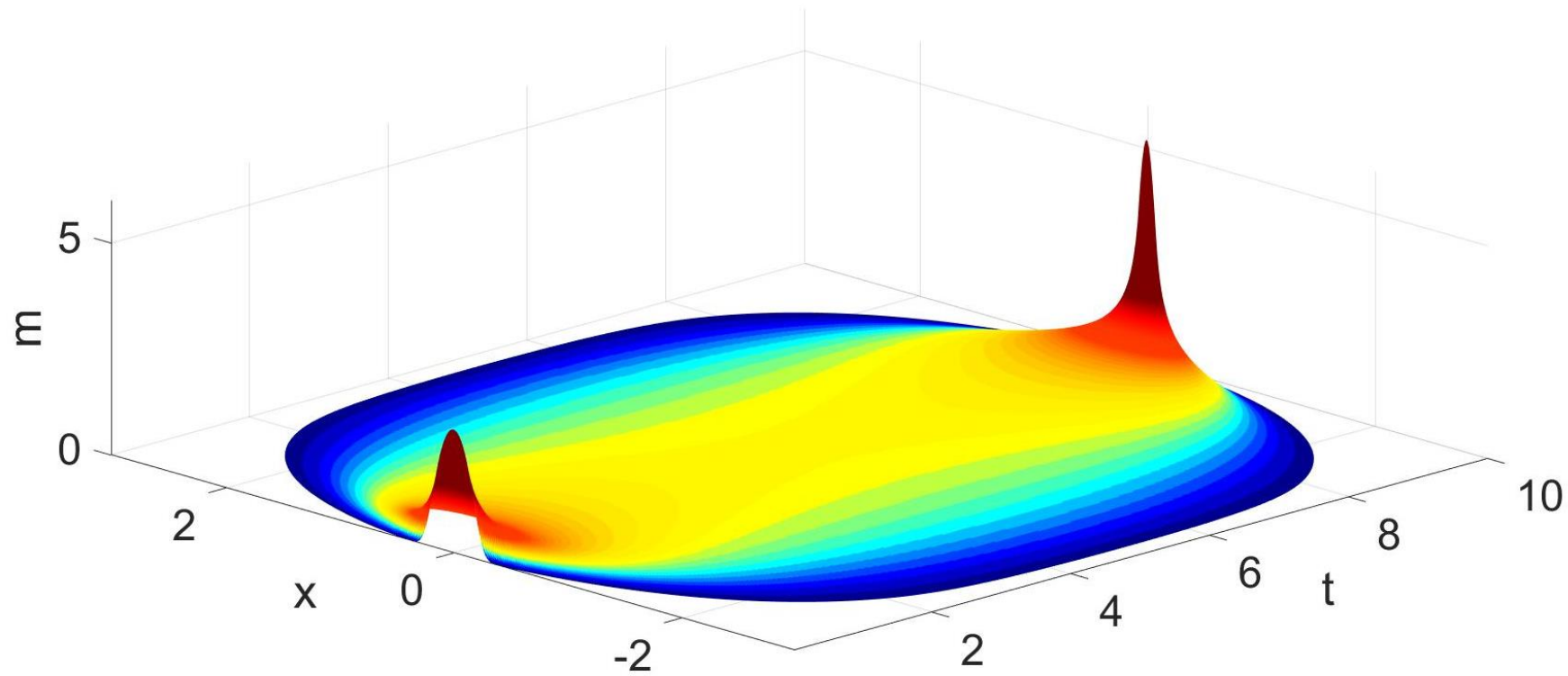


$$\text{for } 0 \ll t \ll T \quad \left\{ \begin{array}{l} m(\mathbf{y}, t) \simeq m_e(\mathbf{y}) \\ u(\mathbf{y}, t) \simeq u_e(\mathbf{y}) + \lambda t \end{array} \right.$$

$$(m_e, u_e, \lambda) \text{ such that } \begin{cases} \lambda - \mathbf{v} \cdot \nabla_{\mathbf{y}} - \frac{1}{2\mu} (\nabla_{\mathbf{y}} u_e)^2 + \frac{\sigma^2}{2} \Delta_{\mathbf{y}} u_e = V[m_e](\mathbf{x}) \\ \nabla_{\mathbf{y}} (m_e (\nabla_{\mathbf{y}} u_e) + \mathbf{v}) - \frac{\sigma^2}{2} \Delta_{\mathbf{y}} m_e = 0 \end{cases}$$

$= -gm_0$

# A simple example



“Search party” toy model:  $V[m](x, t) = \underbrace{gm(x, t)}_{< 0} + \underbrace{U_0(x)}_{\propto -x^2}$

# Non-linear Schrödinger representation

[Bonnemain, Gobron, Ullmo Phys.Lett. A (2020) ; SciPost (2020) ; J. Math. Phys. (2021)]

Introduce two new variables  $\Phi(\mathbf{x}, t)$ ,  $\Gamma(\mathbf{x}, t)$  defined by:

$$u(\mathbf{x}, t) = -\mu\sigma^2 \log(\Phi(\mathbf{x}, t)) \quad m(\mathbf{x}, t) = \Gamma(\mathbf{x}, t)\Phi(\mathbf{x}, t)$$

$$\Rightarrow \begin{cases} \mu\sigma^2 \partial_t \Gamma = \frac{\mu\sigma^4}{2} \Delta_{\mathbf{x}} \Gamma + U_0(\mathbf{x})\Gamma + g m \Gamma \\ -\mu\sigma^2 \partial_t \Phi = \frac{\mu\sigma^4}{2} \Delta_{\mathbf{x}} \Phi + U_0(\mathbf{x})\Phi + g m \Phi \end{cases}$$

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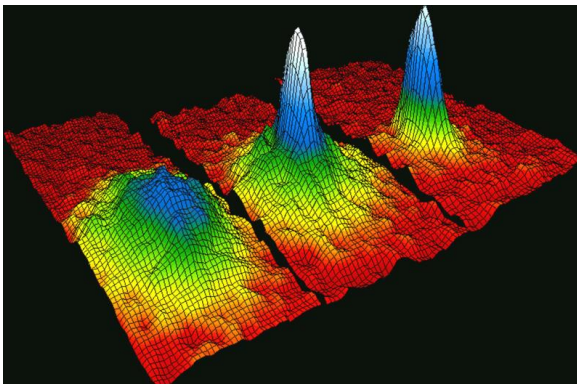
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Rubidium atoms (170 nK)

$$i\hbar \partial_t \Psi = -\frac{\hbar^2}{2\mu} \Delta_{\mathbf{x}} \Psi + U_0(\mathbf{x})\Psi + g|\Psi|^2\Psi$$

Non-Linear Schrödinger

$$(\Psi, \Psi^*, \hbar) \rightarrow (\Phi, \Gamma, i\mu\sigma^2)$$

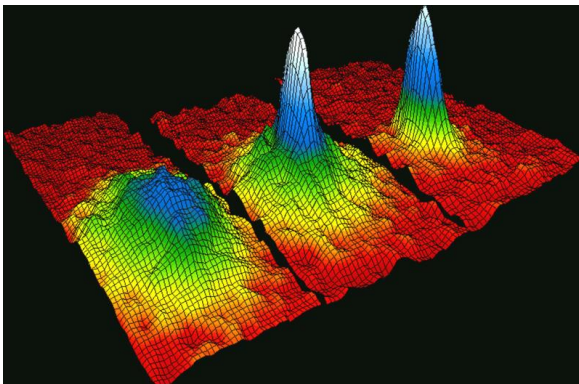
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# Quantum mechanical formalism

[Ullmo, Swiecicki, Gobron (2019)]

- Operators:  $\hat{X} \equiv \mathbf{x}$  ,  $\hat{\Pi} \equiv \mu\sigma^2\nabla_{\mathbf{x}}$  ,  $\hat{O} \equiv f(\hat{\Pi}, \hat{X})$

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e.g.  $\langle \hat{N} \rangle \equiv \int d\mathbf{x} m(\mathbf{x}, t)$  ,  $\langle \hat{X} \rangle \equiv \int d\mathbf{x} m(\mathbf{x}, t) \mathbf{x}$

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$$\begin{cases} -\mu\sigma^2\partial_t\Gamma = \hat{H}\Gamma \\ +\mu\sigma^2\partial_t\Phi = \hat{H}\Phi \end{cases} \Rightarrow \mu\sigma^2\frac{d}{dt}\langle\hat{O}\rangle = \langle\partial_t\hat{O}\rangle + \langle[\hat{O}, \hat{H}]\rangle$$

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e.g.  $\begin{cases} \frac{d}{dt}\langle \hat{X} \rangle = \frac{\langle \hat{\Pi} \rangle}{\mu} \\ \frac{d}{dt}\langle \hat{\Pi} \rangle = \langle -\nabla_{\mathbf{x}}U_0(\hat{X}) \rangle \end{cases} \begin{cases} \frac{d}{dt}\Sigma^2 = \frac{1}{\mu} \left( \langle \hat{X}\hat{\Pi} + \hat{\Pi}\hat{X} \rangle - 2\langle \hat{\Pi} \rangle\langle \hat{X} \rangle \right) \equiv \frac{\Lambda}{\mu} \\ \frac{d}{dt}\Lambda = -2\langle \hat{X}\nabla_{\mathbf{x}}U_0(\hat{X}) \rangle + 2\langle \hat{\Pi}^2 \rangle \end{cases} \quad \Sigma^2 \equiv \langle \hat{X}^2 \rangle - \langle \hat{X} \rangle^2$

# Action and variational approach

- Quadratic MFG are variational systems

$$S[\Gamma, \Phi] \equiv \int_0^T dt \int_{\mathbb{R}} d\mathbf{x} \left[ \frac{\mu\sigma^2}{2} (\Gamma \partial_t \Phi - \Phi \partial_t \Gamma) \right. \\ \left. - \frac{\mu\sigma^4}{2} \nabla \Gamma \cdot \nabla \Phi + \left[ U_0 + \frac{g}{2} \Gamma \Phi \right] \Gamma \Phi \right]$$

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- Noether theorem: Energy is conserved

$$E = \int_{\mathbb{R}} d\mathbf{x} \left[ -\frac{\mu\sigma^4}{2} \nabla\Gamma \cdot \nabla\Phi + \frac{g}{2} (\Gamma\Phi)^2 + U_0\Gamma\Phi \right]$$

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$$\Rightarrow \frac{E^{\text{kin}}}{E^{\text{int}}} \sim \frac{\xi}{L} \quad \xi_{1\text{D}} = \frac{\mu\sigma^4}{|g|} \quad \xi_{2\text{D}} = \sqrt{\frac{\mu\sigma^4}{|g|m_0}}$$



# Solitons and integrability in (1+1)D

[Bonnemain, Gobron, Ullmo (2021)]

- In (1+1)D, if  $U_0 = 0$ : NLS and MFG are integrable.
- Lax connection

$$\partial_t U - \partial_x V + [U, V] = 0$$

$$U = \kappa_\epsilon \begin{pmatrix} \frac{\lambda}{2} & \Phi \\ \Gamma & -\frac{\lambda}{2} \end{pmatrix}, \quad V = \kappa_\epsilon \begin{pmatrix} \kappa_\epsilon \Phi \Gamma & -\partial_x \Phi \\ \partial_x \Gamma & -\kappa_\epsilon \Phi \Gamma \end{pmatrix} - \lambda U$$

- Poisson bracket and infinite hierarchy of conservation laws

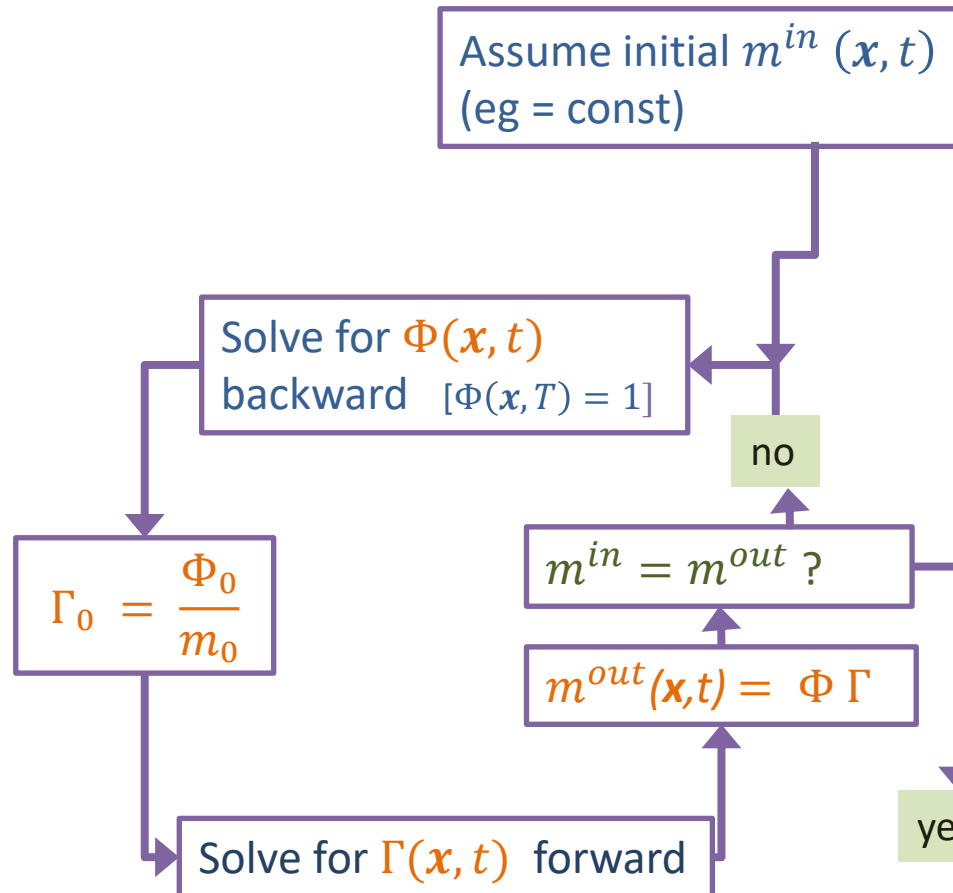
$$\{F, G\} = \int_{\mathbb{R}} \left( \frac{\delta F}{\delta \Gamma} \frac{\delta G}{\delta \Phi} - \frac{\delta F}{\delta \Phi} \frac{\delta G}{\delta \Gamma} \right) dx, \quad \{Q_n, Q_m\} = 0$$

- Soliton solutions

$$\Psi(x, t) = 2b \operatorname{sech} [2b(x + 4at - x_0)] e^{\pm 2[ax + 2(a^2 - b^2)t + \phi_0]}$$

# Back to our model of pedestrian dynamics

## Numerical implementation



Propagation :

$$\begin{cases} -\mu\sigma^2\partial_t\Phi = \frac{\mu\sigma^4}{2}\Delta\Phi + (U_0 + gm^{in})\Phi \\ +\mu\sigma^2\partial_t\Gamma = \frac{\mu\sigma^4}{2}\Delta\Gamma + (U_0 + gm^{in})\Gamma \end{cases}$$

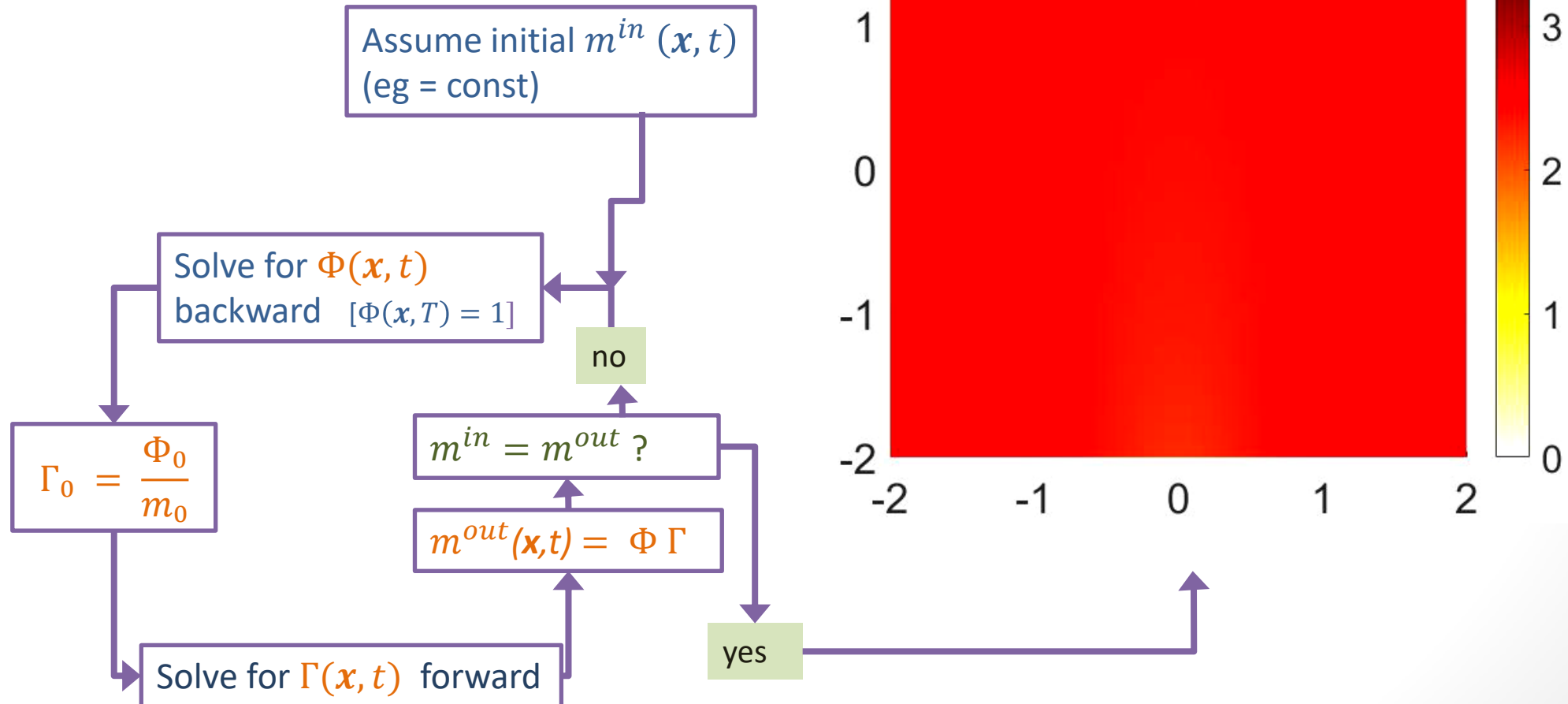
$$[\Phi(T, \cdot) = 1] \quad [\Gamma(0, \cdot) = m_0(\cdot)/\Phi(0, \cdot)]$$

Self consistent equation :

$$m^{out}(\mathbf{x}, t) \equiv \Gamma(\mathbf{x}, t)\Phi(\mathbf{x}, t) = m^{in}(\mathbf{x}, t)$$

# Back to our model of pedestrian dynamics

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# Ergodic vs time dependent

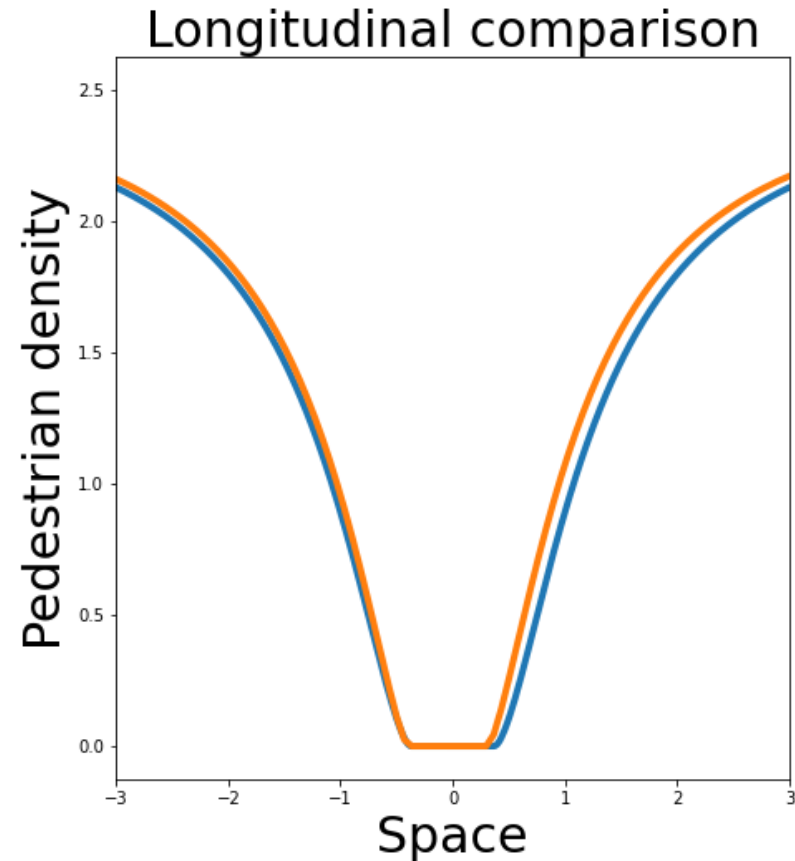
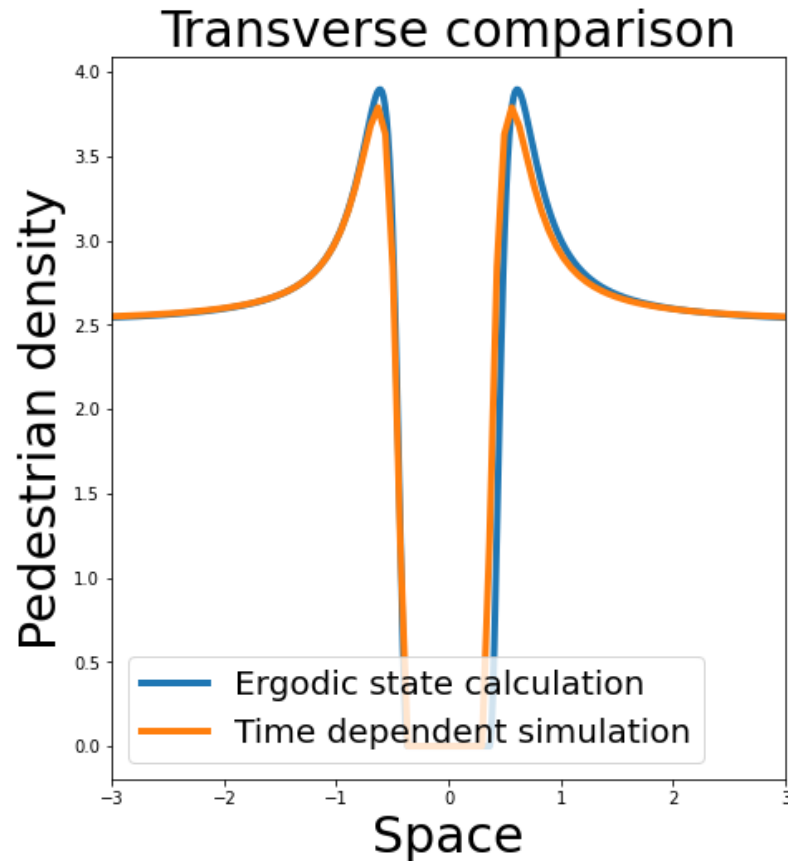
NB : exact symmetry for ergodic case

$$\begin{cases} -gm_0\Gamma_e(\mathbf{y}) = \frac{\mu\sigma^4}{2}\Delta_{\mathbf{y}}\Gamma_e(\mathbf{y}) + \mathbf{v} \cdot \nabla_{\mathbf{y}}\Gamma_e(\mathbf{y}) + U_0(\mathbf{y})\Gamma_e(\mathbf{y}) + g m_e(\mathbf{y})\Gamma_e(\mathbf{y}) \\ -gm_0\Phi_e(\mathbf{y}) = \frac{\mu\sigma^4}{2}\Delta_{\mathbf{y}}\Phi_e(\mathbf{y}) - \mathbf{v} \cdot \nabla_{\mathbf{y}}\Phi_e(\mathbf{y}) + U_0(\mathbf{y})\Phi_e(\mathbf{y}) + g m_e(\mathbf{y})\Phi_e(\mathbf{y}) \end{cases}$$

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# Characteristic length and velocity scales

Intruder

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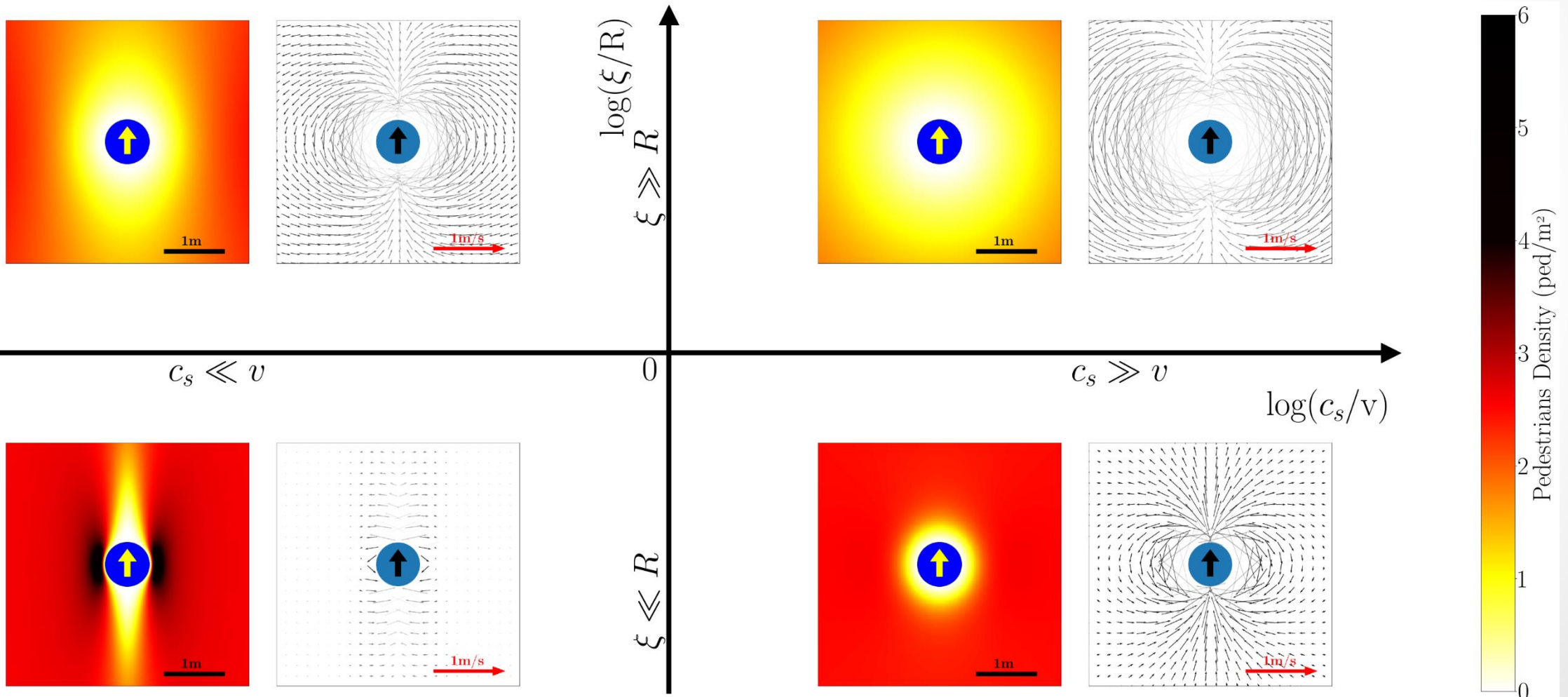
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Up to a scaling factor, solutions of the (ergodic) MFG equations depend only on  $\xi/R$  and  $c_s/v$ .

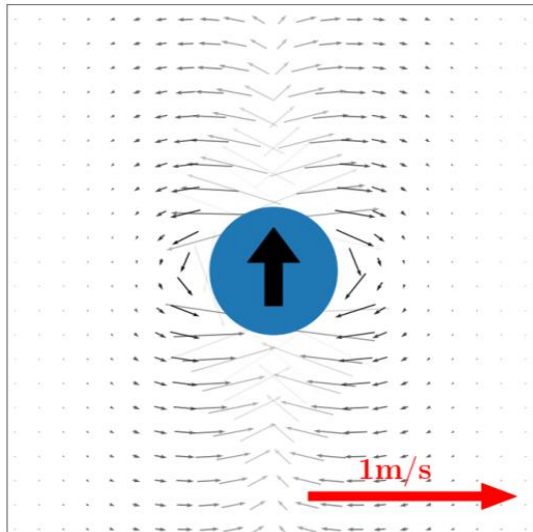
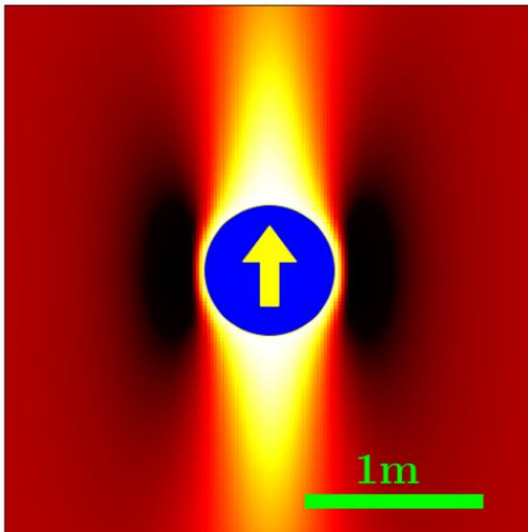
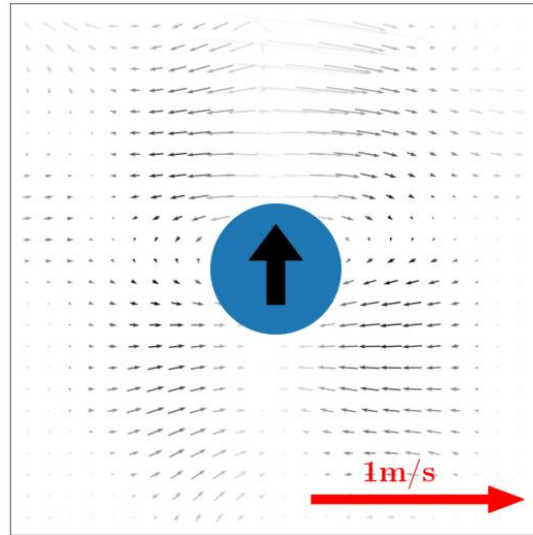
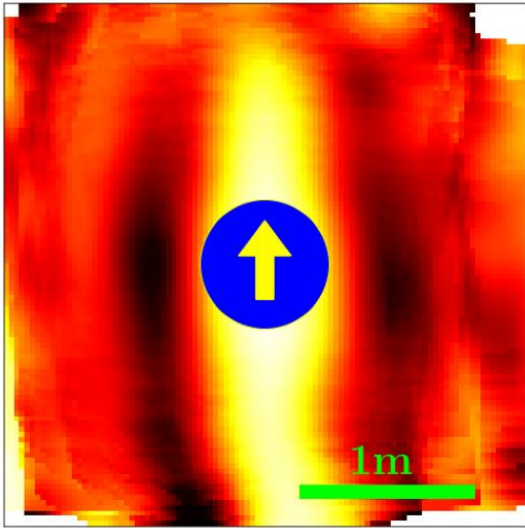


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# Comparison to experiment



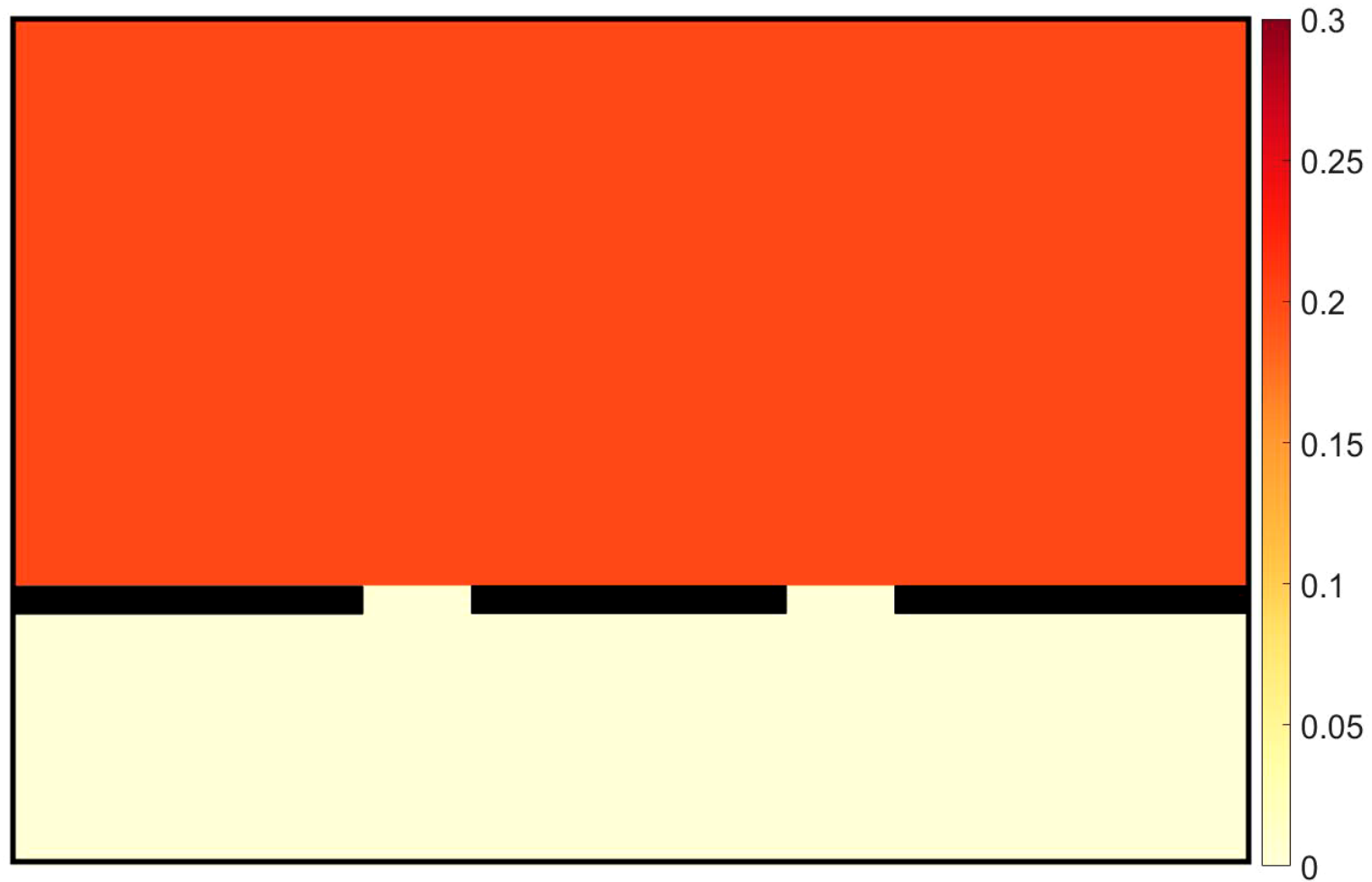
Good qualitative agreement.

Does better than the other models.

Cannot claim quantitative agreement.

# Not limited to obstacle avoidance

[Bonnemain et al. PRE (2023)]



# Discounted Mean-Field Games

[Butano, Appert-Rolland, Ullmo (2023)]

- Add a discount factor

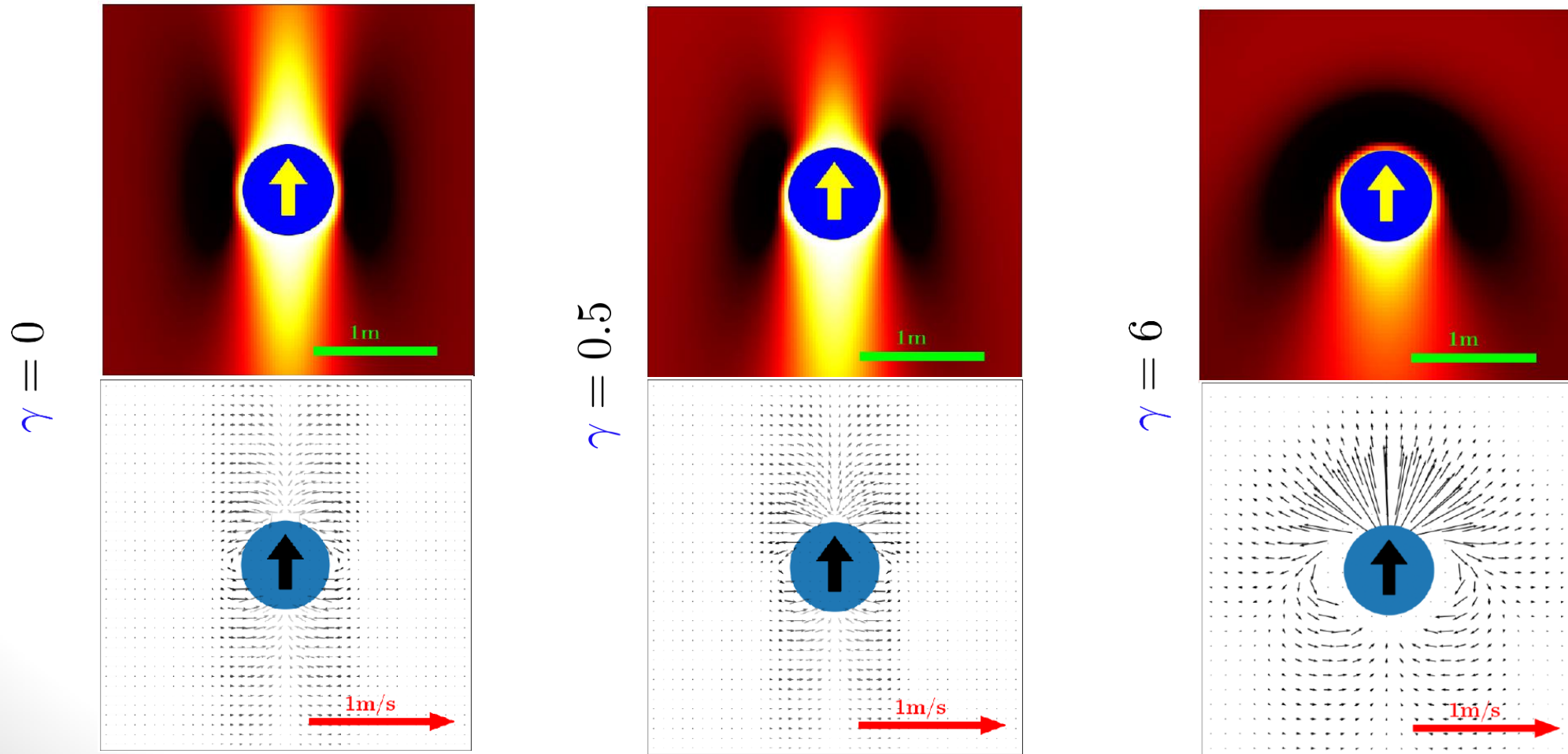
$$\mathbb{E} \left[ \int_t^T \left( \frac{\mu}{2} (\mathbf{a}_s^i)^2 - gm(\mathbf{X}_s^i, s) - U_0(\mathbf{X}_s^i - \mathbf{v}t) \right) e^{\gamma(t-s)} ds + c_T(\mathbf{X}_T^i) e^{\gamma(t-T)} \right]$$

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But...

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And thankfully...

- MFG is versatile.
- Room for improvement (e.g. congestion [Lachapelle, Wolfram (2011)], discount factor, more accurate interaction potential, feedback on the intruder...).